

Quantum Physics 1

Class 13

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Simple Harmonic Oscillator

Review:

- is Functional Vector Space
- is Energy eigenstates of the time independent S.E.

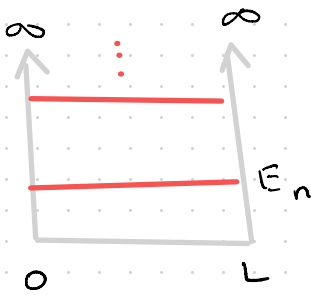
$$\Psi(x), \varphi_n, \dots$$

$$\hat{x}, \hat{p}, \hat{H}, \dots$$

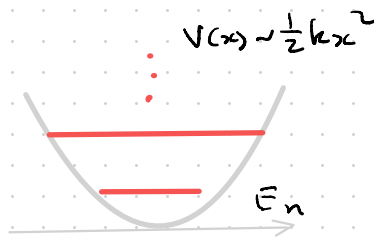
$$\underbrace{-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \varphi_n + V(x) \varphi_n}_{\hat{H}} = E_n \varphi_n(x)$$

$\hat{H} \equiv$ Hamiltonian operator

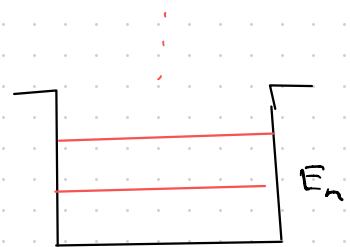
Examples of different eigenstates that manifest from different $V(x)$:



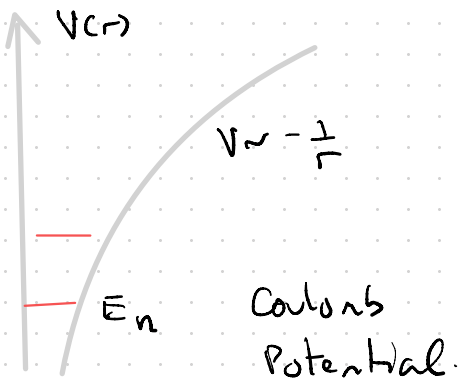
∞ - square well



Harmonic Potential well



Finite square well



- Can write $\Psi(x) = \sum c_n \ell_n(x)$, expansion of eigenstates of the $\hat{H}$. Solutions of time dependent S.E.:

$$\Psi(x,t) = \sum_n c_n \ell_n(x) e^{-i E_n t / \hbar}$$

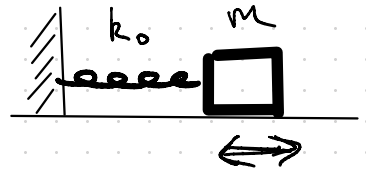
- Can solve fourier coefficients by considering

$$\begin{aligned} \int \ell_n^* \Psi(x,t=0) dx &= \int \sum c_n \ell_n^*(x) \ell_n(x) dx \\ &= \sum c_n \delta_{nn} \\ &= \underline{\underline{c_n}} \end{aligned}$$

Recall: Classical Harmonic Oscillator

$$V = \frac{1}{2} k_0 x^2, \quad k_0 \equiv \text{spring constant}$$

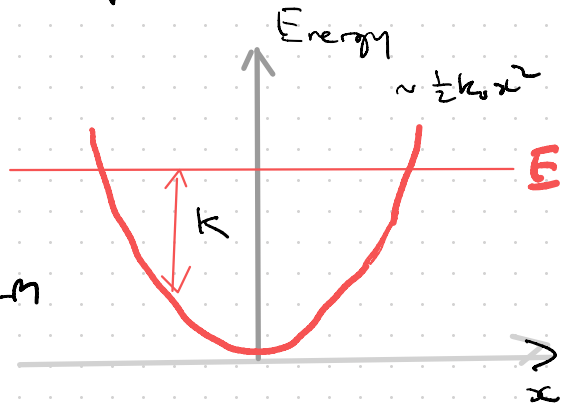
$$F = -\frac{\partial V}{\partial x} = -k_0 x = m \frac{d^2 x}{dt^2}$$



with solutions: $e^{i\omega t}$; $\omega \equiv \sqrt{k_0/m}$

$$\hat{H} = \frac{\hat{p}^2}{2m} + V(x) \Rightarrow$$

$$K = p^2/2m$$

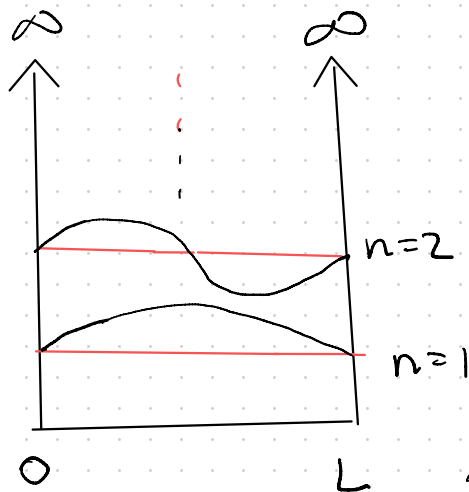


$$\therefore E = \frac{p^2}{2m} + V(x)$$

Quantum Case:

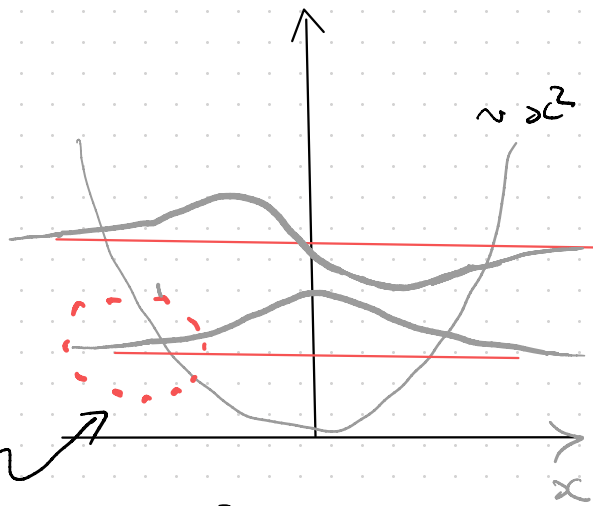
$$\hat{H} = \frac{\hat{p}^2}{2m} + V(\hat{x}) \Rightarrow \frac{-\hbar^2}{2m} \frac{\partial^2 \psi(x)}{\partial x^2} + V(x) \psi(x) = E \psi(x)$$

In-class 13.1, 13.2



∞ - square well

wave penetration



Quantum Harmonic Oscillator (QHO)

Solu:

e^{ikx} for $\left[\frac{p^2}{2m} + V(x) \right] \psi = E \psi$

$\uparrow$ $= 0$ inside, ∞ outside

For QHO $V(x) \sim \frac{1}{2} k x^2$

Eigenfunctions for the QHO

$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi(x) + \frac{1}{2} k_0 x^2 \psi(x) = E \psi(x)$$

with $\psi_n(x) = A_n H_n\left(\sqrt{\frac{m\omega}{\hbar}} x\right) e^{-m\omega x^2/2\hbar}$

↗
Hermite Polynomials

$$H_n(y) = (-1)^n e^{y^2} \frac{d^n}{dy^n} e^{-y^2}$$

$$E_n = \left(n + \frac{1}{2}\right) \hbar \omega ; \quad n = 0, 1, 2, \dots$$

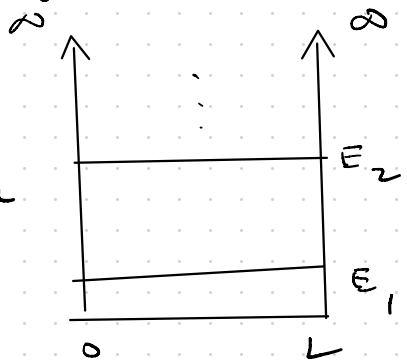
NB: $n=0$ Ground state
with $E_0 = \frac{1}{2} \hbar \omega$

$$\begin{aligned} \Delta E_{n, n+1} &= \hbar \omega ; \quad E_{n+1} - E_n = \left[n+1 + \frac{1}{2}\right] \hbar \omega - \left[n + \frac{1}{2}\right] \hbar \omega \\ &= \hbar \omega \end{aligned}$$

In-class 13.3, 13.4

Compare the Infinite Square Well :

$$\begin{aligned}\Delta E_{n+1,n} &= E_{n+1} - E_n \\ &= \frac{\hbar^2 (n+1)^2 \pi^2}{2mL^2} - \frac{\hbar^2 n^2 \pi^2}{2mL^2} \\ &= \frac{2n\pi^2 \hbar^2}{2mL^2} + \frac{\pi^2 \hbar^2}{2mL^2}\end{aligned}$$

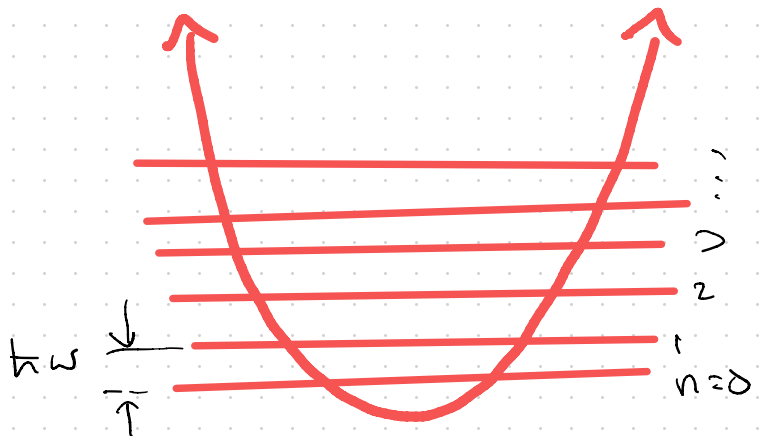


$\propto n$ (dependent on n !)

$\therefore$ if $\Delta E_{n+1,n} \uparrow$ then $n \uparrow$

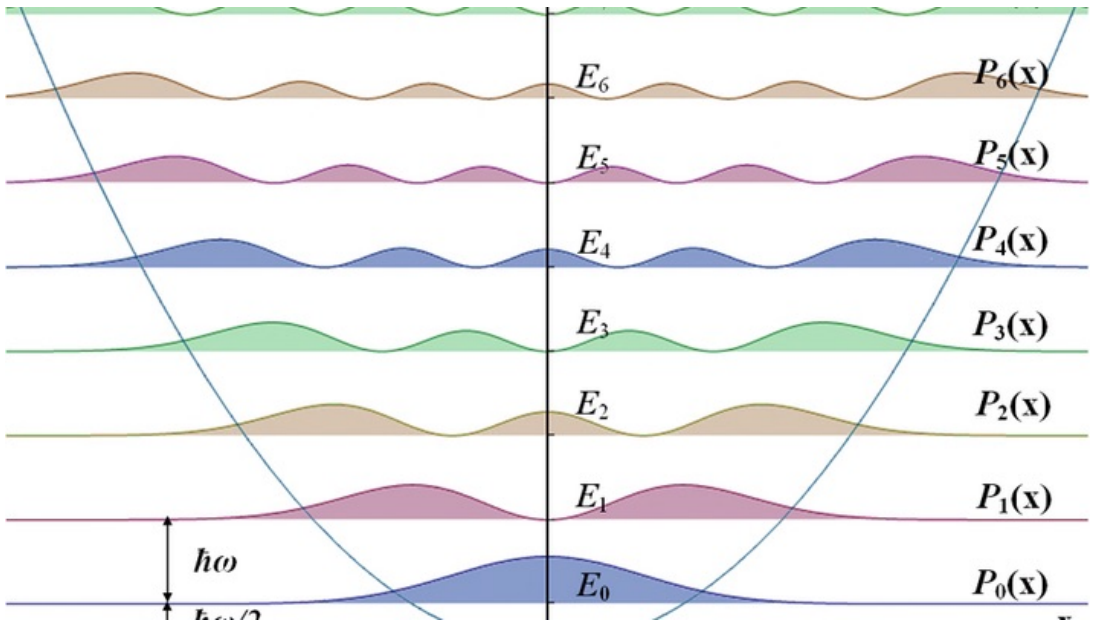
But for QHO: $\Delta E = \hbar\omega$;

equally spaced and independent of n .



Applications :

1. Phonons
 2. Molecular vibrations.
- Any system where can make parabolic approximation at small E, V .



The Science Of Agents Of S.H.I.E.L.D.: What Is Quantum Harmonic Oscillation?

<https://www.forbes.com/sites/chadorzel/2015/10/13/the-science-of-agents-of-s-h-i-e-l-d-what-is-quantum-harmonic-oscillation/#7c2b81c3e7e3>