

Quantum Physics 1

Class 10

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Functional vector space

Last Time :

We solved S.E. using separation of variables

$$\underline{\Psi}(x, t) = \Psi(x) f(t)$$

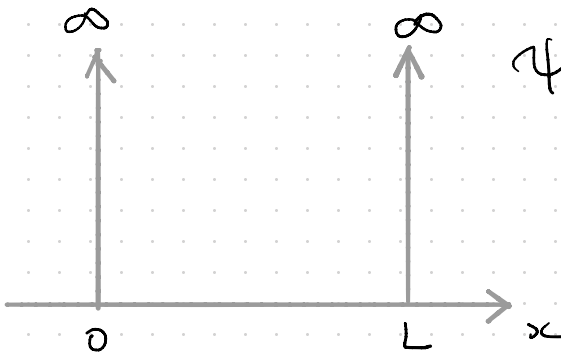
$$\left\{ \begin{array}{l} \frac{\hbar^2}{2m} \frac{\partial^2 \Psi(x)}{\partial x^2} + V(x) \Psi(x) = E \Psi(x) \dots (1) \end{array} \right.$$

$$\frac{\partial f}{\partial t} = -i \frac{E}{\hbar} f(t) \dots (2)$$

$$\Rightarrow f(t) = f(0) e^{-iEt/\hbar}$$

Infinite Square Well :

↗ normaliz. const.



$$\Psi(x) = \sqrt{\frac{2}{L}} \sin(k_n x)$$

$$E_n = \frac{n^2 \pi^2 \hbar^2}{2mL^2}$$

↖ kinetic energy

VIP points:

- (1) Energy is quantized
- (2) Ground state energy > 0
- (3) $\psi^* \psi$ is time independent.

In-class 10.1, 10.2,
10.3

Recall superposition:

$$\Psi(x, t) = \sum c_n \varphi_n(x) e^{-i E_n t / \hbar}$$

at $t=0$ $\Psi(x, 0) = \sum c_n \varphi_n(x)$ ** Fourier Series Expansion
 $\hookrightarrow$ what is c_n ?

$$\Rightarrow \Psi(x, 0) = \sum c_n \varphi_n(x)$$

multiply both sides by φ_m^*

$$\Rightarrow \varphi_m^* \Psi(x, 0) = \sum c_n \varphi_m^* \varphi_n(x)$$

$$\Rightarrow \int \varphi_m^* \Psi(x, 0) dx = \sum c_n \underbrace{\int \varphi_m^* \varphi_n(x) dx}$$

$$\underbrace{\int \varphi_m^* \varphi_n dx}$$

$$\int \sin\left(\frac{n\pi x}{L}\right) \sin\left(\frac{n\pi x}{L}\right) dx$$

$$= \begin{cases} 0 & m \neq n \\ 1 & m = n \end{cases}$$

$$= \delta_{m,n} \text{ Kronecker delta.}$$

$$\Rightarrow \int \varphi_m^* \Psi(x, 0) dx = \sum c_n \delta_{m,n} \\ = c_m$$

$$\therefore c_m = \int \varphi_m^* \Psi(x, 0) dx \equiv \text{Fourier coefficient.}$$

NB : $|c_m|^2 \sim$ probability of a particle being in the m^{th} state.

Now, recall for $t > 0$

$$\Psi(x, t) = \sum_n c_n \varphi_n(x) e^{-i E_n t / \hbar}$$

Consider: $\int_0^L \bar{\Psi}^*(x,t) \bar{\Psi}(x,t) dx = 1$

$$= \int_0^L \sum_n c_n^* \phi_n^*(x) \sum_m c_m \phi_m(x) dx$$

$$= \sum_n \sum_m c_n^* c_m \delta_{m,n}$$

$$= \sum_n |c_n|^2 = 1$$

↪ probability of finding
particle in the
 n^{th} state

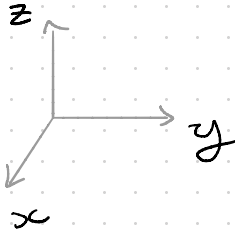
In-class 10.4

Functional Vector Space

Recall Geometrical vector space-

vectors: $\vec{A}, \vec{B}, \dots$

in Euclidean space



$$\begin{aligned}\vec{A} &= A_x \hat{i} + A_y \hat{j} + A_z \hat{k} \\ &= \sum A_i \hat{e}_i ; \hat{e}_1 = \hat{i} \\ &\quad \hat{e}_2 = \hat{j} \\ &\quad \hat{e}_3 = \hat{k}\end{aligned}$$

recall:

$$(i) \hat{e}_i \cdot \hat{e}_j = \delta_{ij}$$

$$\begin{aligned}(ii) A \hat{i} &= \vec{A} \cdot \hat{e}_1 \\ &= A_x\end{aligned}$$

Compare with:

Wavefunction vector space

with objects:

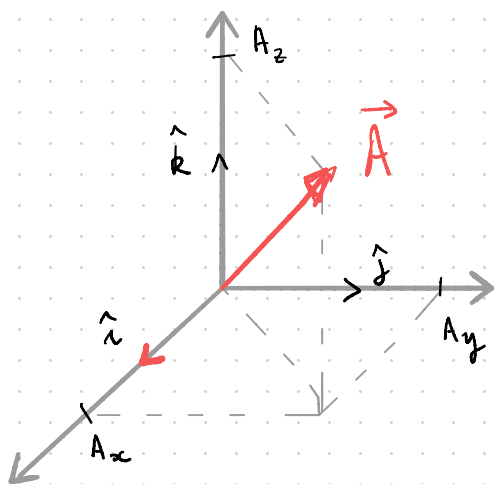
$$\Psi(x, t), \phi_n, \phi_n^*$$

$\phi_n \Rightarrow$ unit vectors
(basis vectors)

$$\Psi(x) = \sum c_n \phi_n(x)$$

$$c_n = \int \phi_n^* \Psi(x) dx$$

Geometrical vector space



Now,

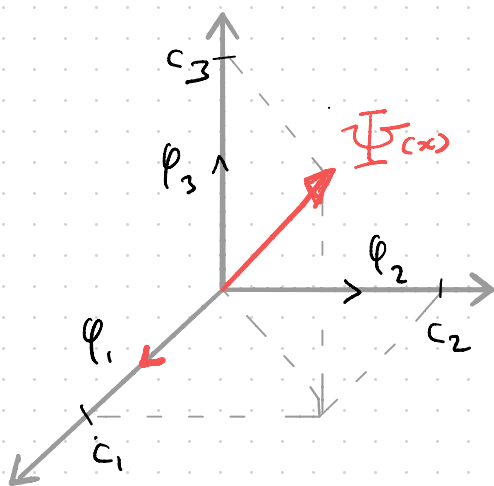
$$\hat{r}_i \rightarrow \varphi_n$$

$$A_i \rightarrow c_n$$

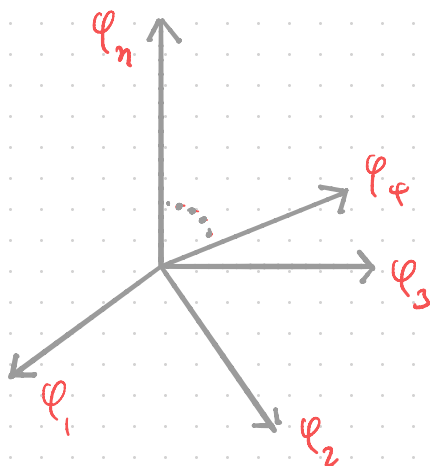
$$\begin{cases} \vec{A} \cdot \hat{r}_i = A_{xi} \\ c_n = \int \varphi_n^*(x) \bar{\Psi}(x) dx \end{cases}$$

$$\begin{cases} \hat{r}_i \cdot \hat{r}_j = \delta_{ij} \\ \int \varphi_m^+ \varphi_n dx = \delta_{m,n} \end{cases}$$

Wavefunction vector space:



in general:
 $\exists$ a High-dimensional space



Inner product:

$$\vec{A} = \sum A_n \hat{n} = A_x \hat{i} + A_y \hat{j} + A_z \hat{k}$$

$$\vec{B} = \sum B_n \hat{n} = B_x \hat{i} + B_y \hat{j} + B_z \hat{k}$$

$$\begin{aligned} \vec{A} \cdot \vec{B} &= \left(\sum_n A_n \hat{n} \right) \cdot \left(\sum_m B_m \hat{n} \right) \\ &= A_x B_x + A_y B_y + A_z B_z \end{aligned}$$

[Dot product]

$$\underline{\Psi}(x) = \sum c_n \phi_n(x)$$

$$\underline{\phi}(x) = \sum D_n \phi_n(x)$$

$$\int \underline{\Psi}^*(x) \underline{\phi}(x)$$

$$= \int \sum_n \sum_m c_n^* \phi_n^*(x) D_m \phi_m(x) dx$$

$$= \sum_n \sum_m c_n^* D_m \delta_{nm}$$

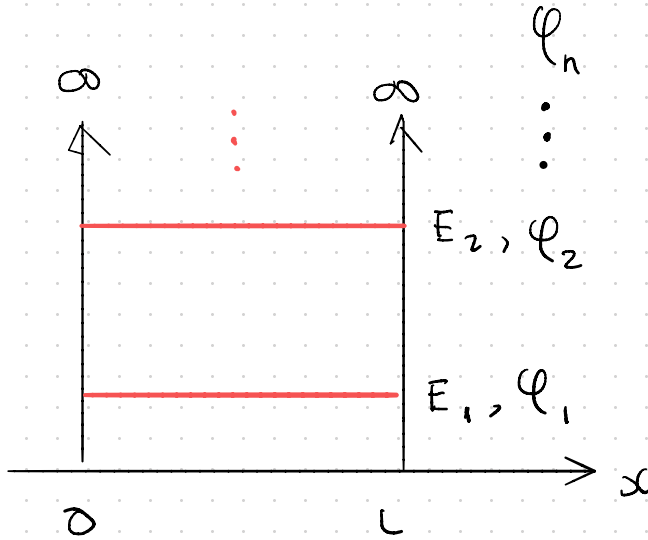
$$= \sum_n c_n^* D_n$$

Now, reconsider the time independent Schrödinger equation for infinite square well.

$$\Rightarrow \phi_n(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right) \rightarrow \text{basis vector}$$

$$\Rightarrow \underbrace{\left[-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x) \right]}_{\equiv \hat{H}} \phi(x) = E \phi(x)$$

$$[\text{Aside: Classically } \hat{H} = \frac{p^2}{2m} + V]$$



where $\varphi_n(x) \equiv$ Eigenfunctions of the Hamiltonian $\hat{H}$

